



Central Spectral Gaps of the Almost Mathieu Operator

I. Krasovsky

Department of Mathematics, Imperial College, London SW7 2AZ, UK. E-mail: i.krasovsky@ic.ac.uk

Received: 20 March 2016 / Accepted: 8 September 2016

Published online: 20 October 2016 – © The Author(s) 2016. This article is published with open access at Springerlink.com

Abstract: We consider the spectrum of the almost Mathieu operator H_α with frequency α and in the case of the critical coupling. Let an irrational α be such that $|\alpha - p_n/q_n| < cq_n^{-\kappa}$, where $p_n/q_n, n = 1, 2, \dots$ are the convergents to α , and c, κ are positive absolute constants, $\kappa < 56$. Assuming certain conditions on the parity of the coefficients of the continued fraction of α , we show that the central gaps of $H_{p_n/q_n}, n = 1, 2, \dots$, are inherited as spectral gaps of H_α of length at least $c'q_n^{-\kappa/2}, c' > 0$.

1. Introduction

Let $H_{\alpha,\theta}$ with $\alpha, \theta \in (0, 1]$ be the self-adjoint operator acting on $l^2(\mathbb{Z})$ as follows:

$$(H_{\alpha,\theta}\phi)(n) = \phi(n-1) + \phi(n+1) + 2\cos 2\pi(\alpha n + \theta)\phi(n), \quad n = \dots, -1, 0, 1, \dots \quad (1.1)$$

This operator is known as the almost Mathieu, Harper, or Azbel-Hofstadter operator. It is a one-dimensional discrete periodic (for α rational) or quasiperiodic (for α irrational) Schrödinger operator which models an electron on the 2-dimensional square lattice in a perpendicular magnetic field. Analysis of the spectrum of $H_{\alpha,\theta}$ (and its natural generalization when the prefactor 2 of cosine, the coupling, is replaced by an arbitrary real number λ) has been a subject of many investigations. In the present paper, we are concerned with the structure of the spectrum of $H_{\alpha,\theta}$ as a set. Denote by $a_j \in \mathbb{Z}_+, j = 1, 2, \dots$ (infinite sequence if α is irrational) the coefficients of the continued fraction of α :

$$\alpha = [a_1, a_2, \dots] = \frac{1}{a_1 + \frac{1}{a_2 + \dots}}.$$

If $\alpha = p/q$ is rational, where p, q are coprime, i.e. $(p, q) = 1$, positive integers, there exists n such that

$$p/q = [a_1, a_2, \dots, a_n] = \frac{1}{a_1 + \frac{1}{a_2 + \dots + \frac{1}{a_n}}}.$$

We denote by $S(\alpha)$ the union of the spectra of $H_{\alpha, \theta}$ over all $\theta \in (0, 1]$. (Note, however, that if α is irrational, the spectrum of $H_{\alpha, \theta}$ does not depend on θ). If $\alpha = p/q$, $S(p/q)$ consists of q bands separated by gaps. As shown by van Mouche [22] and by Choi, Elliott, and Yui [7], all the gaps (with the exception of the centermost gap when q is even) are open. Much effort was expended to prove the conjectures of [1, 4] that if α is irrational, the spectrum is a Cantor set. B ellissard and Simon proved in [6] that the spectrum of the generalized operator mentioned above is a Cantor set for an (unspecified) dense set of pairs (α, λ) in \mathbb{R}^2 . Helffer and S j strand [11] proved the Cantor structure and provided an analysis of gaps in the case when all the coefficients a_j 's of α are sufficiently large. Choi, Elliott, and Yui [7] showed that in the case of $\alpha = p/q$, each open gap is at least of width 8^{-q} (this bound was improved in [3] to $e^{-\varepsilon q}$ with any $\varepsilon > 0$ for q sufficiently large) which, together with a continuity result implies that all admissible gaps are open (in particular, the spectrum is a Cantor set) if α is a Liouvillean number whose convergents p/q satisfy $|\alpha - p/q| < e^{-Cq}$. Last [16] showed that $S(\alpha)$ has Lebesgue measure zero (and hence, since $S(\alpha)$ is closed and known not to contain isolated points, a Cantor set) for all $\alpha = [a_1, a_2, \dots]$ such that the sequence $\{a_j\}_{j=1}^\infty$ is unbounded. The set of such α 's has full measure 1. On the other hand, it was shown by Puig [23] that in the generalized case $\lambda \neq \pm 2, 0$, the spectrum is a Cantor set for α satisfying a Diophantine condition. Finally, Avila and Krikorian [2] completed the proof that the spectrum for $\lambda = 2$ has zero measure, and hence a Cantor set, for all irrational α 's; moreover, the proof of the fact that the spectrum is a Cantor set for all real $\lambda \neq 0$ and irrational α was completed by Avila and Jitomirskaya in [3]. The measure of the spectrum for any irrational α and real λ is $|4 - 2|\lambda||$: in the case $\lambda \neq \pm 2$, proved for a.e. α also in [16] and for all irrationals in [12]. Also available are bounds on the measure of the union of all gaps, see [8, 14, 17]. Furthermore, see [19] for a recent work on the Hausdorff dimension of the spectrum, and [20], on the question of whether all admissible gaps are open.

In order to have a quantitative description of the spectrum, one would like to know if the exponential $e^{-\varepsilon q}$ estimates for the sizes of the individual gaps can be improved at least for some of the gaps.

In this paper we provide a power-law estimate $Cq^{-\kappa}$, $\kappa < 28$, for the widths of *central* gaps of $S(p/q)$, i.e., the gaps around the centermost band (Theorem 3 below), on a parity condition for the coefficients a_k in $p/q = [a_2, a_2, \dots, a_n]$.

From this result we deduce that $S(\alpha)$ has an infinite number of power-law bounded gaps for any irrational $\alpha = [a_1, a_2, \dots]$ admitting a power-law approximation by its convergents $p_n/q_n = [a_1, a_2, \dots, a_n]$ and with a parity condition on a_j 's (Theorem 4 below). These gaps are inherited from the central ones of $S(p_n/q_n)$, $n = 1, 2, \dots$.

First, let $\alpha = p/q$, $(p, q) = 1$. A standard object used for the analysis of $H_{\alpha, \theta}$ is the discriminant

$$\sigma(E) = -\operatorname{tr} \left\{ \begin{pmatrix} E - 2 \cos(2\pi p/q + \pi/2q) & -1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} E - 2 \cos(2\pi 2p/q + \pi/2q) & -1 \\ 1 & 0 \end{pmatrix} \dots \right. \\ \left. \begin{pmatrix} E - 2 \cos(2\pi qp/q + \pi/2q) & -1 \\ 1 & 0 \end{pmatrix} \right\}, \quad (1.2)$$

a polynomial of degree q in E with the property that $S(p/q)$ is the image of $[-4, 4]$ under the inverse of the mapping $\sigma(E)$. The fact that $S(p/q)$ consists of q bands separated by

$q - 1$ open gaps (except for the centermost empty gap for q even) means that all the zeros of $\sigma(E)$ are simple, in all the maxima the value of $\sigma(E)$ is strictly larger than 4, while in all the minima, strictly less than -4 (except for $E = 0$ for q even, where $|\sigma(0)| = 4$ and the derivative $\sigma'(0) = 0$). Note an important fact that $\sigma(E) = (-1)^q \sigma(-E)$, and hence $S(p/q)$ is symmetric w.r.t. $E = 0$.

In what follows, we assume that q is odd. The case of even q can be considered similarly. Let us number the bands from left to right, from $j = -(q-1)/2$ to $j = (q-1)/2$. Let λ_j denote the centers of the bands, i.e. $\sigma(\lambda_j) = 0$. Note that, by the symmetry of $\sigma(E)$, $\lambda_0 = 0$. Let μ_j and η_j denote the edges of the bands, i.e. $|\sigma(\mu_j)| = |\sigma(\eta_j)| = 4$, assigned as follows. If $q = 4k+3$, $k = 0, 1, \dots$, we set $\sigma(\mu_j) = 4$, $\sigma(\eta_j) = -4$ for all j . (In this case the derivative $\sigma'(0) > 0$, as follows from the fact that $\sigma(E) < 0$ for all E sufficiently large.) If $q = 4k+1$, $k = 0, 1, \dots$, we set $\sigma(\mu_j) = -4$, $\sigma(\eta_j) = 4$ for all j . (In this case the derivative $\sigma'(0) < 0$.) Thus, in both cases, the bands are $B_j = [\eta_j, \mu_j]$ for $|j|$ even, and $B_j = [\mu_j, \eta_j]$ for $|j|$ odd.

Let $w_j = \mu_j - \lambda_j$, $w'_j = \lambda_j - \eta_j$ for $|j|$ even, and $w_j = \eta_j - \lambda_j$, $w'_j = \lambda_j - \mu_j$ for $|j|$ odd. Thus, the width of the j 's band is always $w_j + w'_j$. By the symmetry, for the centermost band $B_0 = [\eta_0, \mu_0]$, $w_0 = w'_0$, and in general $w_j = w'_{-j}$.

For any real α , denote the gaps of $S(\alpha)$ by $G_j(\alpha)$ and their length by $\Delta_j(\alpha)$. For $\alpha = p/q$, we order them in the natural way, namely,

$$G_j = (\mu_j, \mu_{j+1}), \quad \Delta_j = \mu_{j+1} - \mu_j, \quad \text{for } |j| \text{ even}, \quad (1.3)$$

$$G_j = (\eta_j, \eta_{j+1}), \quad \Delta_j = \eta_{j+1} - \eta_j, \quad \text{for } |j| \text{ odd}. \quad (1.4)$$

By the symmetry, $\Delta_j = \Delta_{-j-1}$ for $0 \leq j < (q-1)/2$.

In Sect. 2, we prove

Lemma 1 (Comparison of the widths for the gaps and bands). *Let $q \geq 3$ be odd. There hold the inequalities*

$$\Delta_0 > \left(\frac{w_0}{4}\right)^2, \quad \Delta_j > \frac{w_j^2}{4C_0^{2(j+1)}}, \quad 1 \leq j < \frac{q-1}{2}, \quad (1.5)$$

$$\Delta_j > \left(\frac{w_0}{8}\right)^{2j}, \quad 1 \leq j < \frac{q-1}{2}, \quad (1.6)$$

where $C_0 = 1 + 2e/(\sqrt{5} - 1) = 5.398\dots$

Remark. The inequalities of Lemma 1 are better for small j , i.e., for central gaps and bands, which is the case we need below. For large j , note the following estimate which one can deduce using the technique of Last [16]: $\Delta_j > \min\{w_j^2, w'_{j+1}{}^2\}/(4q)$, $0 \leq j < (q-1)/2$.

The inequality (1.6) gives us a lower bound for the width of the j 's gap provided an estimate for the width of the 0's band can be established. Such an estimate is given by

Lemma 2 (Bound for the width of the centermost band). *Let $q \geq 1$, $p/q = p_n/q_n = [a_1, a_2, \dots, a_n]$, where a_1 is odd and a_k , $2 \leq k \leq n$ are even. Then there exist absolute constants $1 < C_1 < 14$ and $1 < C_2 < e^{10}$ such that for the derivative of $\sigma(E)$ at zero*

$$|\sigma'(0)| < C_2 q^{C_1}, \quad (1.7)$$

and half the width of the centermost band of $S(p/q)$

$$w_0 \geq \frac{4}{|\sigma'(0)|} > 4C_2^{-1}q^{-C_1}. \quad (1.8)$$

If, in addition, $q_{k+1} \geq q_k^v$, for some $v > 1$ and all $1 \leq k \leq n-1$, then for any $\varepsilon > 0$ there exists $Q = Q(\varepsilon, v)$ such that if $q > Q$,

$$|\sigma'(0)| < q^{5+\gamma_0+\varepsilon}, \quad w_0 > 4q^{-(5+\gamma_0+\varepsilon)}, \quad (1.9)$$

where γ_0 is Euler's constant.

Remark. The bounds on C_1, C_2 can be somewhat improved.

This lemma is proved in Sect. 3. The inequalities (1.5), (1.6), and especially (1.7) are the main technical results of this paper.

Combination of Lemmata 1 and 2 immediately yields

Theorem 3 (Bound for the widths of the gaps). *Let $q \geq 3$, $p/q = [a_1, a_2, \dots, a_n]$, where a_1 is odd and a_k , $2 \leq k \leq n$, are even. Then, with C_k , $k = 1, 2$, from Lemma 2, the width of the j 's gap of $S(p/q)$ is*

$$\Delta_0 > \left(\frac{1}{C_2 q^{C_1}} \right)^2, \quad \Delta_j > \left(\frac{1}{2C_2 q^{C_1}} \right)^{2j}, \quad 1 \leq j < \frac{q-1}{2}. \quad (1.10)$$

Remark. The improvements for large q on the additional condition $q_{k+1} > Cq_k^v$ are obvious from (1.9).

A consequence of this is the following theorem proved in Sect. 4.

Theorem 4 *There exists an absolute $C_3 > 0$ such that the following holds. Let $\alpha = [a_1, a_2, \dots] \in (0, 1)$ be an irrational such that a_1 is odd, a_k , $k \geq 2$, are even, and such that*

$$\left| \alpha - \frac{p_n}{q_n} \right| < \frac{1}{C_3 q_n^\kappa}, \quad \kappa = 4C_1, \quad (1.11)$$

for all $p_n/q_n = [a_1, a_2, \dots, a_n]$, $n = 1, 2, \dots$, where C_1 is the constant from Lemma 2. Then

- (a) *The interior of the centermost band B_0 of $S(p_n/q_n)$ contains the centermost band and the closures of the gaps G_0, G_{-1} of $S(p_{n+1}/q_{n+1})$, $n = 1, 2, \dots$*
- (b) *There exist distinct gaps $G_{n,j}(\alpha)$, $n = 1, 2, \dots$, $j = 1, 2$, of $S(\alpha)$, such that the intersections $G_{n,1}(\alpha) \cap G_{-1}(p_n/q_n)$, $G_{n,2}(\alpha) \cap G_0(p_n/q_n)$, $n = 1, 2, \dots$ are non-empty and the length of the gap $G_{n,j}(\alpha)$*

$$\begin{aligned} \Delta_{n,j}(\alpha) &= |G_{n,j}(\alpha)| \geq |G_{n,j}(\alpha) \cap G_{j-2}(p_n/q_n)| \\ &> \frac{1}{C_4 q_n^{\kappa/2}}, \quad n = 1, 2, \dots, \quad j = 1, 2, \end{aligned} \quad (1.12)$$

for some absolute $C_4 > 0$, where $|A|$ denotes the Lebesgue measure of A .

- (c) *Let $\varepsilon > 0$, replace C_3 by 2, and set $\kappa = 4(5 + \gamma_0 + \varepsilon)$ in (1.11). Then there exists $n_0 = n_0(\varepsilon)$ such that (a) and (b) hold for all $n = n_0, n_0 + 1, \dots$ (instead of $n = 1, 2, \dots$) with C_4 replaced by 2, and with $\kappa/2$ in (1.12) replaced by $2(5 + \gamma_0) + \varepsilon$.*

Remarks. (1) The statements (a), (b) of the theorem hold a fortiori for $\kappa = 4 \cdot 14 = 56$ and for any larger κ . It is easy to provide explicit examples of irrationals satisfying the conditions of Theorem 4: take $\kappa = 56$, any odd a_1 , and even a_{n+1} such that $a_{n+1} > C_3 q_n^{\kappa-2}$, $n \geq 1$. Indeed, in this case,

$$\left| \alpha - \frac{p_n}{q_n} \right| < \frac{1}{q_n q_{n+1}} < \frac{1}{a_{n+1} q_n^2} < \frac{1}{C_3 q_n^\kappa}.$$

- (2) Note that the parity condition on a_j 's implies, in particular, that all q_n 's are odd. This condition can be relaxed in all our statements. For example, we can allow a finite number of a_j 's to be odd at the expense of excluding some $G(p_n/q_n)$'s from the statement of Theorem 4 and worsening the bound on C_1 . Note that in Lemma 2 we need q to be odd in order to use the estimate (1.7) on $\sigma'(0)$ to obtain (1.8). One could obtain a bound on w_0 for even q by providing an estimate on the second derivative $\sigma''(0)$ in this case: for q even $\sigma'(0) = 0$. The parity condition we assume in this paper allows the best estimates and simplest proofs.
- (3) In Theorem 4, we only use Theorem 3 for $j = 0$, i.e., for the 2 centermost gaps. One can extend the result of Theorem 4, with appropriate changes, to more than 2 (at least a finite number) of central gaps of $S(p_n/q_n)$.
- (4) We can take $C_3 = 4^2 60^2 C_2^4$, $C_4 = 2C_2^2$, in terms of the constant C_2 from Lemma 2.
- (5) The statement (a) of the theorem holds already for $\kappa = 2C_1$.

2. Proof of Lemma 1

Assume that $q = 4k + 3$, $k = 0, 1, \dots$ (A proof in the case $q = 4k + 1$ is almost identical.) Let

$$s = \frac{q-1}{2}.$$

In our notation, we can write

$$\begin{aligned} \sigma(E) &= \prod_{k=-s}^s (E - \lambda_k), & \sigma(E) - 4 &= \prod_{k=-s}^s (E - \mu_k), \\ \sigma(E) + 4 &= \prod_{k=-s}^s (E - \eta_k). \end{aligned} \quad (2.1)$$

Setting in the last 2 equations $E = \lambda_j$, we obtain the useful identities

$$4 = \prod_{k=-s}^s |\lambda_j - \mu_k|, \quad 4 = \prod_{k=-s}^s |\lambda_j - \eta_k|, \quad -s \leq j \leq s. \quad (2.2)$$

Fix $0 \leq j \leq s$ (by the symmetry of the spectrum, it is sufficient to consider only nonnegative j). It was shown by Choi, Elliott, and Yui [7] that

$$\prod_{k \neq j} |\mu_j - \mu_k| \geq 1, \quad \prod_{k \neq j} |\eta_j - \eta_k| \geq 1. \quad (2.3)$$

For simplicity of notation, we assume from now on that $j < s - 1$: the extension to $j = s - 1$ is obvious. Let $j \geq 0$ be even. By the first inequality in (2.3), we can write

$$\begin{aligned} 1 \leq |\sigma'(\mu_j)| &= \prod_{k \neq j} |\mu_j - \mu_k| \\ &= |\mu_j - \mu_{j+1}| \frac{\prod_{k=-s}^s |\lambda_j - \mu_k|}{|\lambda_j - \mu_j| |\lambda_j - \mu_{j+1}|} \prod_{k=-s}^{j-1} \left| 1 + \frac{\mu_j - \lambda_j}{\lambda_j - \mu_k} \right| \prod_{k=j+2}^s \left| 1 - \frac{\mu_j - \lambda_j}{\mu_k - \lambda_j} \right|. \end{aligned} \quad (2.4)$$

According to our notation, $\mu_{j+1} - \mu_j = \Delta_j$, $\mu_j - \lambda_j = w_j$, $\mu_{j+1} - \lambda_j = w_j + \Delta_j$. Recalling the first identity in (2.2) and rearranging the last product in (2.4), we continue (2.4) as follows

$$= \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \frac{\prod_{k=-s}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right|}{\prod_{k=j+2}^s \left| 1 + \frac{w_j}{\mu_k - \mu_j} \right|} < \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \frac{\prod_{k=-s}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right|}{\prod_{k=j+2}^s \left| 1 + \frac{w_j}{\mu_k - \lambda_j} \right|}, \quad (2.5)$$

because $\mu_j > \lambda_j$.

Now note that, by the symmetry of the spectrum,

$$|\mu_{j+\ell} - \lambda_j| < |\mu_{-j-\ell-1} - \lambda_j|, \quad \ell = 2, 3, \dots \quad (2.6)$$

Therefore, the r.h.s. of (2.5) is

$$< \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \prod_{k=-j-2}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right| \frac{1}{1 + \frac{w_j}{\mu_s - \lambda_j}}. \quad (2.7)$$

In the case $j = 0$, we now use the symmetry

$$w_0 = w'_0 < \lambda_0 - \mu_{-k}, \quad k \geq 1, \quad (2.8)$$

to obtain from (2.7)

$$1 < \frac{16\Delta_0}{w_0(w_0 + \Delta_0)} < \frac{16\Delta_0}{w_0^2},$$

which gives the first inequality in (1.5).

In general, however, we need to compare w_j and w'_j to estimate (2.7). According to equations (3.11), (3.12) of Last [16],

$$w_j, w'_j < e\ell_j, \quad \ell_j = \frac{4}{|\sigma'(\lambda_j)|}, \quad (2.9)$$

and further, by equations (3.27), (3.28) of [16],

$$\frac{\sqrt{5}-1}{2}\ell_j < w_j, w'_j. \quad (2.10)$$

(In fact, more is shown in [16]: for each pair of widths w_j, w'_j , at least one of them is larger than ℓ_j .)

Therefore,

$$w_j < c_1 w'_j, \quad c_1 = \frac{2e}{\sqrt{5}-1}, \quad 0 \leq j \leq s. \quad (2.11)$$

Furthermore, it is obvious that

$$\frac{w'_j}{\lambda_j - \mu_k} < 1, \quad k = -j-2, \dots, j-1. \quad (2.12)$$

Therefore, we have for the product in (2.7):

$$\prod_{k=-j-2}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \mu_k} \right| < (1 + c_1)^{2(j+1)}, \quad (2.13)$$

and since $\mu_s - \lambda_j > 0$, (2.7) finally gives

$$1 < \frac{4\Delta_j}{w_j(w_j + \Delta_j)} (1 + c_1)^{2(j+1)}, \quad (2.14)$$

from which the inequality (1.5) with even j easily follows.

Remark. Last's equation (2.9) together with the Last–Wilkinson formula [16, 18]

$$\sum_{j=-s}^s |\sigma'(\lambda_j)|^{-1} = 1/q \quad (2.15)$$

implies [16] that the measure of the spectrum $S(p/q)$ is at most $8e/q$ and that for any j ,

$$w_j < 4e/q. \quad (2.16)$$

Now consider j odd, $0 < j < s$. Using the second inequalities in (2.3) and (2.2), we obtain similarly to (2.5),

$$1 < \frac{4\Delta_j}{w_j(w_j + \Delta_j)} \frac{\prod_{k=-s}^{j-1} \left| 1 + \frac{w_j}{\lambda_j - \eta_k} \right|}{\prod_{k=j+2}^s \left| 1 + \frac{w_j}{\eta_k - \lambda_j} \right|}, \quad (2.17)$$

and since

$$|\eta_{j+\ell} - \lambda_j| < |\eta_{-j-\ell-1} - \lambda_j|, \quad \ell = 2, 3, \dots, \quad (2.18)$$

we obtain the inequality (1.5) for j odd in a similar way.

Let again j be even, $0 < j < s-1$. In order to compare Δ_j with the width of the centermost band and, thus, obtain (1.6), we write instead of (2.4) the following:

$$1 \leq |\sigma'(\mu_j)| = \prod_{k \neq j} |\mu_j - \mu_k| \quad (2.19)$$

$$= |\mu_j - \mu_{j+1}| \frac{\prod_{k=-s}^s |\lambda_0 - \mu_k|}{|\lambda_0 - \mu_j| |\lambda_0 - \mu_{j+1}|} \prod_{k=-s}^{j-1} \left| 1 + \frac{\mu_j - \lambda_0}{\lambda_0 - \mu_k} \right| \prod_{k=j+2}^s \left| 1 - \frac{\mu_j - \lambda_0}{\mu_k - \lambda_0} \right|. \quad (2.20)$$

Proceeding in a similar way as before, and using the inequalities

$$|\mu_{j+\ell} - \lambda_0| < |\mu_{-j-\ell-1} - \lambda_0|, \quad \ell = 2, 3, \dots, \quad (2.21)$$

we obtain

$$\begin{aligned} 1 &< \frac{4\Delta_j}{|\lambda_0 - \mu_j||\lambda_0 - \mu_{j+1}|} \prod_{k=-j-2}^{j-1} \left| 1 + \frac{\mu_j - \lambda_0}{\lambda_0 - \mu_k} \right| \\ &< \frac{16\Delta_j}{|\lambda_0 - \mu_j||\lambda_0 - \mu_{j+1}|} \prod_{k=-j}^{j-1} \left| 1 + \frac{\mu_j - \lambda_0}{\lambda_0 - \mu_k} \right| \\ &= \frac{16\Delta_j}{|\lambda_0 - \mu_j||\lambda_0 - \mu_{j+1}|} \left| \frac{\mu_j - \mu_0}{\lambda_0 - \mu_0} \right| \left| \frac{\mu_j - \mu_1}{\lambda_0 - \mu_1} \right| \prod_{\substack{k=-j \\ k \neq 0,1}}^{j-1} \left| \frac{\mu_j - \mu_k}{\lambda_0 - \mu_k} \right| \\ &< \frac{16\Delta_j}{|\lambda_0 - \mu_0||\lambda_0 - \mu_1|} \prod_{\substack{k=-j \\ k \neq 0,1}}^{j-1} \left| \frac{\mu_j - \mu_k}{\lambda_0 - \mu_k} \right|, \end{aligned} \quad (2.22)$$

and since (note that $S(\alpha) \in [-4, 4]$)

$$\left| \frac{\mu_j - \mu_k}{\lambda_0 - \mu_k} \right| < \frac{8}{w_0},$$

we obtain

$$1 < \frac{\Delta_j}{4} \left(\frac{8}{w_0} \right)^{2j}, \quad (2.23)$$

which gives an inequality slightly better than (1.6) for j even. Finally, we establish (1.6) for j odd by starting (instead of (2.19)) with the inequality $1 \leq |\sigma'(\eta_{j+1})| = \prod_{k \neq j+1} |\eta_{j+1} - \eta_k|$ and arguing similarly.

Remark. Using the Last estimate (2.9)

$$\frac{4e}{w_j} > |\sigma'(\lambda_j)| = \prod_{k \neq j} |\lambda_j - \lambda_k|, \quad (2.24)$$

one can establish, in a way similar to the argument above, inequalities of the type

$$\Delta_j + w'_{j+1} > \frac{w_j}{C^j}, \quad (2.25)$$

with some absolute constant $C > 0$.

□

3. Proof of Lemma 2

As noted in a remark following Theorem 4, the parity conditions imposed on p/q in Lemma 2 imply, in particular, that q is odd. It follows from the symmetry of the discriminant $\sigma(E) = -\sigma(-E)$ in this case that the maximum of the absolute value of the derivative $\sigma'(E)$ in the $j = 0$ band is at $E = 0$. Therefore,

$$w_0 \geq \frac{4}{|\sigma'(0)|} \quad (3.1)$$

(with the equality only for $q = 1$), and hence, in order to prove Lemma 2, it remains to obtain the inequalities (1.7) and (1.9).

If $q = 1$, we have $\sigma(E) = -E$, and the result is trivial. Assume now that q is any (even or odd) integer larger than 1. We start with the following representation of $\sigma(E)$ in terms of a $q \times q$ Jacobi matrix with the zero main diagonal:

$$\sigma(E) = \det(\hat{H} - EI), \quad (3.2)$$

where I is the identity matrix, and \hat{H} is a $q \times q$ matrix \hat{H}_{jk} , $j, k = 1, \dots, q$, where

$$\hat{H}_{jj+1} = \hat{H}_{j+1j} = 2 \sin\left(\pi \frac{p}{q} j\right), \quad j = 1, \dots, q-1, \quad (3.3)$$

and the rest of the matrix elements are zero. For a proof, see e.g. the appendix of [15]. (This is related to a matrix representation for the almost Mathieu operator corresponding to the chiral gauge of the magnetic field potential, noticed by several authors [13, 21, 25].) The absence of the main diagonal in \hat{H} allows us to obtain a simple expression for the derivative $\sigma'(E)$ at $E = 0$. If q is even, it is easily seen that $\sigma'(E) = 0$. If q is odd, we denote $s = (q-1)/2$ and immediately obtain from (3.2) (henceforth we set $\prod_{j=a}^b \equiv 1$ and $\sum_{j=a}^b \equiv 0$ if $a > b$):

$$\sigma'(0) = (-1)^s \sum_{k=0}^s \left[\prod_{j=1}^k 2 \sin \frac{\pi p}{q} (2j-1) \prod_{j=k+1}^s 2 \sin \frac{\pi p}{q} 2j \right]^2. \quad (3.4)$$

From now on, we assume that $q \geq 3$ is odd unless stated otherwise.

Remark. Using the identity $\prod_{j=1}^{(q-1)/2} 2 \sin \frac{\pi p}{q} 2j = q$, we can represent (3.4) in the form

$$\sigma'(0) = (-1)^s q \left(1 + \sum_{k=1}^s \prod_{j=1}^k \frac{\sin^2 \frac{\pi p}{q} (2j-1)}{\sin^2 \frac{\pi p}{q} 2j} \right), \quad (3.5)$$

which exhibits the fact that $|\sigma'(0)| > q$. This is in accordance with the Last–Wilkinson formula (2.15).

Thus we have

$$|\sigma'(0)| = \sum_{k=0}^s \exp\{L_k\}, \quad (3.6)$$

where

$$\frac{1}{2}L_k = \sum_{j=1}^k \ln \left| 2 \sin \frac{\pi p}{q} (2j-1) \right| + \sum_{j=k+1}^s \ln \left| 2 \sin \frac{\pi p}{q} 2j \right| \quad (3.7)$$

$$= \sum_{j=-s+k+1}^k \ln \left| 2 \sin \frac{\pi p}{q} (2j-1) \right| = \sum_{j=1}^s \ln \left| 2 \sin \frac{\pi p}{q} 2(j+k) \right|. \quad (3.8)$$

Here we changed the summation variable $j = s + j'$ in the second sum in (3.7) to obtain the first equation in (3.8), and then changed the variable $j = k - s + j'$ to obtain the final equation in (3.8).

We will now analyze L_k . Using the Fourier expansion, we can write

$$L_k = -2 \sum_{j=1}^s \sum_{n=1}^{\infty} \frac{1}{n} \cos 4n \frac{\pi p}{q} (j+k). \quad (3.9)$$

Representing n in the form $n = qm + \ell$, where $\ell = 1, 2, \dots, q-1$ for $m = 0$, and $\ell = 0, 1, \dots, q-1$ for $m = 1, 2, \dots$, we have

$$L_k = - \sum_{m=1}^{\infty} \left(\frac{q-1}{qm} - \frac{1}{q} \sum_{\ell=1}^{q-1} \frac{1}{m + \ell/q} F(\ell, k) \right) + S_k, \quad (3.10)$$

where

$$S_k = \sum_{\ell=1}^{q-1} \frac{1}{\ell} F(\ell, k), \quad (3.11)$$

and

$$\begin{aligned} F(\ell, k) &= -2 \sum_{j=1}^s \cos 4\ell \frac{\pi p}{q} (j+k) \\ &= -2\Re \sum_{j=1}^s \exp\{4\pi i \frac{p}{q} (j+k)\ell\} = \frac{\cos \pi \frac{p}{q} \ell (4k+1)}{\cos \pi \frac{p}{q} \ell}. \end{aligned} \quad (3.12)$$

Note that since $1 < j+k \leq q-1$, $(p, q) = 1$, and q is odd, we have that $(2j+2k)p \not\equiv 0 \pmod{q}$.

For $0 \leq k \leq s$,

$$\sum_{\ell=1}^{q-1} F(\ell, k) = -2\Re \left(\sum_{j=1}^s \sum_{\ell=1}^{q-1} \exp\{4\pi i \frac{p}{q} (j+k)\ell\} + 1 - 1 \right) = q-1. \quad (3.13)$$

We will use this fact later on.

Recall that

$$\sum_{m=1}^M \frac{1}{m} = \ln M + \gamma_0 + o(1), \quad M \rightarrow \infty, \quad (3.14)$$

where $\gamma_0 = 0.5772 \dots$ is Euler's constant. Recall also Euler's ψ -function

$$\psi(x+1) = \frac{\Gamma'(x+1)}{\Gamma(x+1)} = -\gamma_0 - \sum_{m=1}^{\infty} \left(\frac{1}{m+x} - \frac{1}{m} \right), \quad x \geq 0. \quad (3.15)$$

The function $\psi(x)$ continues to a meromorphic function in the complex plane with first-order poles at nonpositive integers $x = 0, -1, -2, \dots$. The function $\psi(x)$ satisfies the equation

$$\psi(x+1) = \psi(x) + \frac{1}{x}.$$

Expressions (3.14), (3.15) imply

$$\begin{aligned} \sum_{m=1}^M \frac{1}{m+x} &= \sum_{m=1}^M \left(\frac{1}{m+x} - \frac{1}{m} \right) + \sum_{m=1}^M \frac{1}{m} \\ &= \ln M - \psi(x+1) + o(1), \quad M \rightarrow \infty, \end{aligned} \quad (3.16)$$

uniformly for $x \in [0, 1]$, in particular. We now rewrite (3.10) in the form

$$L_k = - \lim_{M \rightarrow \infty} \left[\frac{q-1}{q} \sum_{m=1}^M \frac{1}{m} - \frac{1}{q} \sum_{\ell=1}^{q-1} F(\ell, k) \sum_{m=1}^M \frac{1}{m + \ell/q} \right] + S_k. \quad (3.17)$$

Substituting here (3.14), (3.16), and using (3.13), we finally obtain

$$L_k = -\frac{q-1}{q} \gamma_0 - \frac{1}{q} \sum_{\ell=1}^{q-1} F(\ell, k) \psi(1 + \ell/q) + S_k. \quad (3.18)$$

We will now provide an upper bound for the absolute values of sums in this expression. First, note that the derivative $\psi'(x) \geq 0, x \in [1, 2]$, and $\psi(1) = -\gamma_0, \psi(2) = \psi(1) + 1 = 1 - \gamma_0$. Therefore,

$$\max_{x \in [1, 2]} |\psi(x)| = \max\{|\psi(1)|, |\psi(2)|\} = \gamma_0.$$

Thus, for the first sum in the r.h.s. of (3.18) we have

$$\begin{aligned} \left| \frac{1}{q} \sum_{\ell=1}^{q-1} F(\ell, k) \psi(1 + \ell/q) \right| &\leq \frac{\gamma_0}{q} \sum_{\ell=1}^{q-1} |F(\ell, k)| = \frac{\gamma_0}{q} \sum_{\ell=1}^{q-1} \frac{1}{|\cos(\pi p \ell/q)|} \\ &= \frac{4\gamma_0}{q} \sum_{m=0}^{s-1} \frac{1}{|1 - e^{i(2m+1)\pi/q}|} < \gamma_0 \sum_{m=0}^{s-1} \frac{1}{m + 1/2} \\ &< \gamma_0 \left(2 + \int_0^{s-1} \frac{dx}{x + 1/2} \right) < \gamma_0 (\ln q + 2), \quad q \geq 3. \end{aligned} \quad (3.19)$$

We will need a more subtle estimate for

$$S_k = \sum_{m=1}^{q-1} \frac{1}{m} F(m, k) = \sum_{m=1}^{q-1} \frac{1}{m} \frac{\cos \pi \frac{p}{q} m (4k+1)}{\cos \pi \frac{p}{q} m} \quad (3.20)$$

in the r.h.s. of (3.18). We follow a method of Hardy and Littlewood [9] (see also [24]). It relies on a recursive application of a suitably constructed contour integral.

For $q \geq 3$ odd, $(p, q) = 1$, let

$$I(p/q, \gamma) = -2 \int_{\Gamma_q} \frac{e^{(1+p/q)z}}{(1+e^{z p/q})(1-e^z)} \frac{e^{-\gamma z}}{z} dz, \quad \frac{1}{2} \frac{p}{q} \leq \gamma \leq 1 + \frac{1}{2} \frac{p}{q}, \quad (3.21)$$

where the contour Γ_q are the 2 direct lines parallel to the real axis given by: (1) $\pi i/2 + x$, $x \in \mathbb{R}$, oriented from $-\infty$ to $+\infty$; (2) $2\pi i(q-1/4) + x$, $x \in \mathbb{R}$, oriented from $+\infty$ to $-\infty$. Note that the choice of γ in (3.21) ensures that the integral converges both at $+\infty$ and $-\infty$.

Now again for $q \geq 3$ odd, $(p, q) = 1$, let

$$S(p/q, \gamma) = \sum_{m=1}^{q-1} \frac{e^{\pi i \frac{p}{q} m - 2\pi i \gamma m}}{m \cos(\pi \frac{p}{q} m)}. \quad (3.22)$$

Note that the sum (3.20)

$$S_k = \Re S(p/q, -2kp/q), \quad (3.23)$$

and that the denominators in (3.22) are nonzero.

We will also need the following auxiliary sum, $(p, q) = 1$,

$$T(p/q, \gamma, \delta) = 2 \sum_{n=1}^q \frac{(-1)^\delta e^{2\pi i \frac{p}{q} (n-\frac{1}{2})}}{1 - (-1)^\delta e^{2\pi i \frac{p}{q} (n-\frac{1}{2})}} \frac{e^{-2\pi i \gamma (n-\frac{1}{2})}}{n - \frac{1}{2}}, \quad \delta = 0, 1, \quad (3.24)$$

where we assume that p is odd if $\delta = 0$ and that p and q have opposite parities if $\delta = 1$. These conditions imply that $p(2n-1) \neq q(2m-\delta)$, $m, n \in \mathbb{Z}$, and therefore the denominators in (3.24) are nonzero.

With this notation we have

Lemma 5. *Let q be odd, $(p, q) = 1$, $p > 0$, $q > 1$. Let p/q have the following continued fraction:*

$$\frac{p}{q} = \frac{1}{a + \frac{p'}{q'}}, \quad p' \geq 0, \quad q' > p'. \quad (3.25)$$

Then

$$I(p/q, \gamma) = S(p/q, \gamma) - (-1)^{\varepsilon'} T(p'/q', \gamma', a \bmod 2), \quad \gamma' = \frac{q}{p} \gamma \pmod{1}, \quad (3.26)$$

where $\varepsilon' = 0$ if $\gamma' = \frac{q}{p} \gamma \pmod{2}$, and $\varepsilon' = 1$ otherwise.

Moreover, there holds the bound

$$|I(p/q, \gamma)| < 4 \ln \frac{q}{p} + \frac{5}{e\pi p} + \beta, \quad \beta = 4(e^{-1} + \operatorname{arcsinh}(4/\pi)) = 5.719 \dots \quad (3.27)$$

Remarks.(1) One can take any γ' satisfying the congruence in (3.26).

(2) The bound in (3.27) can be somewhat decreased by improving (3.33), (3.34) below.

Similar can be achieved in (3.38) below.

Proof. Consider the integral $I_\Gamma(p/q, \gamma)$, which has the same integrand as in (3.21), but with integration over some contour Γ . Let $\Gamma_{q,\xi}$ be the following quadrangle traversed in the positive direction: $\Gamma_{q,\xi} = \cup_{j=1}^4 \Gamma_{q,\xi}^{(j)}$, $\xi > 0$, where $\Gamma_{q,\xi}^{(1)} = [\pi i/2 - \xi, \pi i/2 + \xi]$, $\Gamma_{q,\xi}^{(2)} = [\pi i/2 + \xi, 2\pi i(q - 1/4) + \xi]$, $\Gamma_{q,\xi}^{(3)} = [2\pi i(q - 1/4) - \xi, 2\pi i(q - 1/4) + \xi]$, $\Gamma_{q,\xi}^{(4)} = [2\pi i(q - 1/4) - \xi, \pi i/2 - \xi]$.

Recalling the conditions on γ in (3.21), we first note that on the vertical segments, for some constants C which may depend on p, q ,

$$|I_{\Gamma_{q,\xi}^{(2)}}(p/q, \gamma)| \leq C \frac{e^{-\gamma\xi}}{\xi} \leq C \frac{e^{-\frac{p}{2q}\xi}}{\xi} \rightarrow 0, \quad \text{as } \xi \rightarrow \infty, \quad (3.28)$$

$$|I_{\Gamma_{q,\xi}^{(4)}}(p/q, \gamma)| \leq C \frac{e^{-(1+p/q-\gamma)\xi}}{\xi} \leq C \frac{e^{-\frac{p}{2q}\xi}}{\xi} \rightarrow 0, \quad \text{as } \xi \rightarrow \infty. \quad (3.29)$$

and we conclude that

$$I(p/q, \gamma) = \lim_{\xi \rightarrow \infty} I_{\Gamma_{q,\xi}}(p/q, \gamma). \quad (3.30)$$

On the other hand, $I_{\Gamma_{q,\xi}}(p/q, \gamma)$ is given by the sum of residues inside the contour. Clearly, the integrand has poles there at the points:

$$z_m = 2\pi i m, \quad m = 1, \dots, q-1 \quad (3.31)$$

$$\tilde{z}_n = 2\pi i \frac{q}{p}(n-1/2), \quad n = 1, \dots, p. \quad (3.32)$$

Note that for all these m, n , $z_m \neq \tilde{z}_n$ because $m \neq \frac{q}{p} \frac{2n-1}{2}$ as q is odd. Hence we conclude that all the poles inside $\Gamma_{q,\xi}$ are simple. Computing the residues and using the facts that

$$\frac{q}{p} = a + \frac{p'}{q'}, \quad q' = p,$$

we obtain (3.26) by (3.30). Note that the conditions on p', q' in $T(p'/q')$ are fulfilled since q is odd.

Now, in order to obtain the inequality (3.27), we evaluate the integral along the contour Γ_q . On the lower part of it,

$$\begin{aligned} |I_{\frac{\pi i}{2} + \mathbb{R}}(p/q, \gamma)| &\leq 2 \int_0^\infty \frac{e^{-\gamma x} + e^{-(1+\frac{p}{q}-\gamma)x}}{(1 + 2e^{-\frac{p}{q}x} \cos \frac{\pi p}{2q} + e^{-2\frac{p}{q}x})^{1/2} (1 + e^{-2x})^{1/2}} \frac{dx}{(x^2 + \frac{\pi^2}{4})^{1/2}} \\ &< 4 \int_0^\infty \frac{e^{-\frac{p}{2q}x}}{(x^2 + \frac{\pi^2}{4})^{1/2}} dx = 4 \int_0^\infty \frac{e^{-u}}{(u^2 + (\frac{\pi p}{4q})^2)^{1/2}} du. \end{aligned} \quad (3.33)$$

Separating the final integral into 2 parts, one along $(0, 1)$ and another along $(1, \infty)$, we can continue (3.33) as follows:

$$\begin{aligned}
 &< 4 \left(\int_0^1 \frac{du}{(u^2 + (\frac{\pi p}{4q})^2)^{1/2}} + \frac{1}{(1 + (\frac{\pi p}{4q})^2)^{1/2}} \int_1^\infty e^{-u} du \right) \\
 &= 4 \left(-\ln \frac{\pi p}{4q} + \ln \left[1 + \sqrt{1 + \left(\frac{\pi p}{4q} \right)^2} \right] + \frac{e^{-1}}{(1 + (\frac{\pi p}{4q})^2)^{1/2}} \right) \\
 &< 4 \left(\ln \frac{q}{p} + \ln \left[\frac{4}{\pi} + \sqrt{\left(\frac{4}{\pi} \right)^2 + 1} \right] + e^{-1} \right). \tag{3.34}
 \end{aligned}$$

Similarly, we obtain (recall that $q > 1$)

$$\begin{aligned}
 |I_{2\pi i(q - \frac{1}{4}) + \mathbb{R}}(p/q, \gamma)| &\leq 4 \int_0^\infty \frac{e^{-u}}{(u^2 + (\pi(p - \frac{p}{4q}))^2)^{1/2}} dx \\
 &< \frac{4e^{-1}}{\pi(p - \frac{p}{4q})} < \frac{4e^{-1}}{\pi p} \frac{1}{1 - \frac{1}{4q}} < \frac{5}{e\pi p}. \tag{3.35}
 \end{aligned}$$

The sum of (3.34) and (3.35) gives (3.27), and thus we finish the proof of Lemma 5. \square

We will need another similar lemma.

Lemma 6. *Let $(p, q) = 1$, $p > 0$, $q > 1$,*

$$\begin{aligned}
 J(p/q, \gamma, \delta) &= 2 \int_{\Gamma_q} \frac{(-1)^\delta e^{(1+p/q)z}}{(1 - (-1)^\delta e^{zp/q})(1 + e^z)} \frac{e^{-\gamma z}}{z} dz, \\
 \frac{1}{2} \frac{p}{q} &\leq \gamma \leq 1 + \frac{1}{2} \frac{p}{q}, \tag{3.36}
 \end{aligned}$$

where $\delta = \{0, 1\}$, and Γ_q is the same contour as in (3.21). Assume that p is odd if $\delta = 0$, and either $(p - \text{even}, q - \text{odd})$, or $(p - \text{odd}, q - \text{even})$ if $\delta = 1$. Let p/q have the continued fraction (3.25). Then

$$\begin{aligned}
 &J(p/q, \gamma, \delta) \\
 &= T(p/q, \gamma, \delta) + \begin{cases} -S(p'/q', \gamma'), & \text{if } \delta = 0 \\ (-1)^{\varepsilon'} T(p'/q', \gamma', a + 1 \bmod 2), & \text{if } \delta = 1 \end{cases}, \quad \gamma' = \frac{q}{p} \gamma \pmod{1}, \tag{3.37}
 \end{aligned}$$

where $\varepsilon' = 0$ if $\gamma' = \frac{q}{p} \gamma \pmod{2}$, and $\varepsilon' = 1$ otherwise.

Moreover, there holds the bound with β from (3.27)

$$\begin{aligned}
 |J(p/q, \gamma, \delta)| &< A_\delta \left(4 \ln \frac{q}{p} + \frac{5}{e\pi p} + \beta \right), \\
 A_\delta &= \begin{cases} (1 - \cos^2 \pi \frac{p}{q})^{-1/2}, & \text{if } \delta = 0 \\ 1, & \text{if } \delta = 1 \end{cases} \tag{3.38}
 \end{aligned}$$

Proof. We argue as in the proof of Lemma 5. The poles of the integrand in (3.36) inside $\Gamma_{q,\xi}(p/q, \gamma)$ are:

$$z_m = 2\pi i(m - 1/2), \quad m = 1, \dots, q, \quad (3.39)$$

$$\tilde{z}_n = 2\pi i \frac{q}{p}(n - \delta/2), \quad (3.40)$$

and $n = 1, \dots, p - 1$ if $\delta = 0$, while $n = 1, \dots, p$ if $\delta = 1$. Our assumptions on the parity of p, q immediately imply that all $z_n \neq \tilde{z}_m$ and hence all the poles inside $\Gamma_{q,\xi}(p/q, \gamma)$ are simple. Computing the residues we obtain (3.37).

Denote by $J_\Gamma(p/q, \gamma, \delta)$ the integral which has the same integrand as in (3.36), but with integration over some contour Γ . As in the previous proof, consider now an estimate for the integral along the lower part of Γ_q :

$$\begin{aligned} |J_{\frac{\pi i}{2} + \mathbb{R}}(p/q, \gamma)| &\leq 2 \int_0^\infty \frac{e^{-\gamma x} + e^{-(1+\frac{p}{q}-\gamma)x}}{(1 - 2(-1)^\delta e^{-\frac{p}{q}x} \cos \frac{\pi p}{2q} + e^{-2\frac{p}{q}x})^{1/2} (1 + e^{-2x})^{1/2} (x^2 + \frac{\pi^2}{4})^{1/2}} dx \\ &< 4A_\delta \int_0^\infty \frac{e^{-\frac{p}{2q}x}}{(x^2 + \frac{\pi^2}{4})^{1/2}} dx, \end{aligned} \quad (3.41)$$

where A_δ is given in (3.38). The rest of the argument is very similar to that in the proof of Lemma 5. \square

Lemma 7 (recurrence). *Let $q > 1$, $\frac{p}{q} = \frac{p_n}{q_n} = [a_1, \dots, a_n]$ and denote $t_j = [a_j, \dots, a_n]$, $j = 1, \dots, n$. Assume that a_1 is odd, and a_j are even for $j = 2, \dots, n$. Fix $0 \leq k \leq (q - 1)/2$. Then*

$$S(p/q, \gamma_1) = I(t_1, \gamma_1) + \sum_{j=2}^n (-1)^{\varepsilon_j} J(t_j, \gamma_j, 1) + 2(-1)^{\varepsilon_n} e^{-i\pi\gamma_n a_n} \quad (3.42)$$

where $\varepsilon_j = \{0, 1\}$ and

$$\gamma_1 = -2kt_1 + k_1, \quad \gamma_j = k_{j-1}t_j + k_j, \quad j = 2, \dots, n, \quad (3.43)$$

with the sequence $k_j \in \mathbb{Z}$, $j = 1, \dots, n$, chosen so that $\frac{1}{2}t_j \leq \gamma_j \leq 1 + \frac{1}{2}t_j$.

Furthermore, for this p/q , there holds the following bound for the sum (3.20):

$$|S_k| = \left| \sum_{m=1}^{q-1} \frac{1}{m} \frac{\cos \pi \frac{p}{q} m(4k+1)}{\cos \pi \frac{p}{q} m} \right| < \left(4 + \frac{\beta}{\ln 2} \right) \ln q + 9. \quad (3.44)$$

If, in addition, $q_{k+1} \geq q_k^v$, for some $v > 1$ and all $1 \leq k \leq n - 1$, then for any $\varepsilon > 0$ there exist $Q = Q(\varepsilon, v)$ such that if $q > Q$,

$$|S_k| < (4 + \varepsilon) \ln q. \quad (3.45)$$

Proof. First, note that $t_1 = p/q$ and $t_j = 1/(a_j + t_{j+1})$, $j = 1, \dots, n-1$, $t_n = 1/a_n$. Now note a simple fact that the conditions $a_1 - \text{odd}$, $a_2, \dots, a_n - \text{even}$ ensure that q is odd and all the fractions t_j for $j = 2, \dots, n$ are either $\frac{\text{odd}}{\text{even}}$ or $\frac{\text{even}}{\text{odd}}$. We now choose k_1 so that $\gamma_1 = -2kt_1 + k_1$ satisfies $\frac{1}{2}t_1 \leq \gamma_1 \leq 1 + \frac{1}{2}t_1$. Applying Lemma 5 to $S(t_1, \gamma_1)$, we obtain

$$S(t_1, \gamma_1) = I(t_1, \gamma_1) + (-1)^\varepsilon T(t_2, \gamma_2, a_1 \bmod 2) \quad (3.46)$$

We can and will choose γ_2 so that $\frac{1}{2}t_2 \leq \gamma_2 \leq 1 + \frac{1}{2}t_2$ by picking the appropriate k'_2 and setting

$$\begin{aligned} \gamma_2 &= \frac{q}{p}\gamma_1 + k'_2 \\ &= \frac{q}{p}(-2kt_1 + k_1) + k'_2 \\ &= -2k + \frac{q}{p}k_1 + k'_2 = -2k + (a_1 + t_2)k_1 + k'_2 = t_2k_1 + k_2, \end{aligned} \quad (3.47)$$

where $k_2 = -2k + a_1k_1 + k'_2$. The constant ε is determined by γ_2 as described in Lemma 5.

Since $a_1 \bmod 2 = 1$, we can now apply Lemma 6 with $\delta = 1$ to T in the r.h.s. of (3.46). This gives

$$S(t_1, \gamma_1) = I(t_1, \gamma_1) + (-1)^\varepsilon [J(t_2, \gamma_2, 1) - (-1)^{\varepsilon'} T(t_3, \gamma_3, a_2 + 1 \bmod 2)]. \quad (3.48)$$

with

$$\gamma_3 = \frac{1}{t_2}\gamma_2 + k'_3 = k_1 + (a_3 + t_3)k_2 + k'_3 = t_3k_2 + k_3.$$

chosen so that $\frac{1}{2}t_3 \leq \gamma_3 \leq 1 + \frac{1}{2}t_3$.

Note that according to our assumption $a_j + 1 \bmod 2 = 1$, $j = 2, \dots, n$, so that we can continue applying Lemma 6 with $\delta = 1$ in (3.48) recursively. At the final step, revisiting the residue calculations for Lemma 6 gives:

$$T(t_n, \gamma_n, 1) = J(t_n, \gamma_n, 1) + 2e^{-i\pi\gamma_n a_n}, \quad (3.49)$$

which proves (3.42).

Now using the bounds (3.27) and (3.38) for $\delta = 1$, we write

$$\begin{aligned} |S_k| &< |S(p/q, \gamma_1)| \leq |I(t_1, \gamma_1)| + \sum_{j=2}^n |J(t_j, \gamma_j, 1)| + 2 \\ &< 4 \ln \frac{1}{t_1 t_2 \cdots t_n} + \frac{5}{e\pi} \sum_{j=1}^n \frac{1}{p'_j} + n\beta + 2, \end{aligned} \quad (3.50)$$

where $t_j = [a_j, \dots, a_n] = \frac{p'_j}{q'_j}$.

Recall that we denote $q = q_n$, $p = p_n$. Observe that the following recurrence [10] with $t_{n+1} = 0$

$$q_n + t_{n+1}q_{n-1} = (a_n + t_{n+1})q_{n-1} + q_{n-2} = \frac{1}{t_n}(q_{n-1} + t_n q_{n-2})$$

gives

$$\frac{1}{t_1 t_2 \cdots t_n} = q_n. \quad (3.51)$$

Note that this equation holds for any continued fraction.

Furthermore, the recurrence $t_{j-1} = 1/(a_{j-1} + t_j)$ implies

$$p'_{j-1} = q'_j, \quad q'_{j-1} = a_{j-1} p'_{j-1} + p'_j,$$

and so

$$p'_{j-2} = a_{j-1} p'_{j-1} + p'_j,$$

which, with the initial conditions $p'_n = 1$, $p'_{n-1} = a_n$, and our assumption that all a_j , $j = 2, \dots, n$, are even, gives

$$p'_j \geq a_{j+1} \cdots a_n \geq 2^{n-j}, \quad j = 1, \dots, n. \quad (3.52)$$

Thus we have

$$\sum_{j=1}^n \frac{1}{p'_j} < 2. \quad (3.53)$$

Finally, since

$$q_n = a_n q_{n-1} + q_{n-2} \geq a_n q_{n-1} \geq \cdots \geq a_n a_{n-1} \cdots a_1 \geq 2^{n-1},$$

we have

$$n \leq \frac{\ln q_n}{\ln 2} + 1. \quad (3.54)$$

(In fact, as is well known, a slightly worse bound on n holds for any continued fraction).

Using (3.51), (3.53), (3.54) in (3.50), we obtain (3.44).

To obtain (3.45) note first the following.

Lemma 8. *Let $v > 1$, $n \geq 2$, $p_k/q_k = [a_1, a_2, \dots, a_k]$, and such that $q_{k+1} \geq q_k^v$, $1 \leq k \leq n-1$, $q_2 \geq 3$. Then*

$$n \leq \frac{1}{\ln v} \ln \frac{\ln q_n}{\ln 3} + 2$$

Proof. We have

$$q_n \geq q_{n-1}^v \geq q_{n-2}^{v^2} \geq \cdots \geq q_2^{v^{n-2}} \geq 3^{v^{n-2}}, \quad n \geq 2,$$

from which the result follows. \square

Now using Lemma 8 instead of (3.54) in (3.50) and choosing Q sufficiently large, we obtain (3.45) if $q > Q$, and finish the proof of Lemma 7. \square

Bringing together (3.6), (3.18), (3.19), and (3.44), yields the bound

$$|\sigma'(0)| < e^{-\frac{2}{3}\gamma_0} \frac{2}{3} q \cdot q^{\gamma_0+4+\beta/\ln 2} e^{2\gamma_0+9} < q^{\gamma_0+5+\beta/\ln 2} \frac{2}{3} e^{9+4\gamma_0/3}, \quad q \geq 3. \quad (3.55)$$

Since $\gamma_0 + 5 + \beta/\ln 2 = 13.8 \dots < 14$, $\frac{2}{3} e^{9+4\gamma_0/3} < e^{10}$, and $\sigma(E) = -E$ if $q = 1$, we obtain (1.7). Finally, using (3.45) instead of (3.44), we obtain (1.9). This finishes the proof of Lemma 2. \square

4. Proof of Theorem 4

Note first that since for any irrational α

$$\frac{1}{2q_n q_{n+1}} < \frac{1}{q_n(q_n + q_{n+1})} < \left| \alpha - \frac{p_n}{q_n} \right|,$$

we obtain that for any α satisfying the conditions of Theorem 4,

$$\frac{C_3}{2} q_n^{\kappa-1} < q_{n+1}, \quad n = 1, 2, \dots \quad (4.1)$$

We will denote $\mu_j(p/q)$, $w_j(p/q)$, etc, the values μ_j , w_j , etc, for the spectrum $S(p/q)$. Fix $n \geq 1$. Let E_2 , E_0 be the right edges of the centermost bands in $S(p_n/q_n)$, $S(p_{n+1}/q_{n+1})$, respectively. By Lemma 2,

$$E_2 = \mu_0(p_n/q_n) = w_0(p_n/q_n) > \frac{4}{C_2 q_n^{C_1}}. \quad (4.2)$$

On the other hand by (2.16) and (4.1), we have

$$E_0 = \mu_0(p_{n+1}/q_{n+1}) = w_0(p_{n+1}/q_{n+1}) < \frac{4e}{q_{n+1}} < \frac{8e}{C_3 q_n^{\kappa-1}}. \quad (4.3)$$

We will now show that $G_0(p_{n+1}/q_{n+1}) \subset (E_0, E_2 - \epsilon)$ for a suitably chosen C_3 .

Recall a continuity property found by Avron, Van Mouche, Simon [5]: if $E \in S(\beta)$, there is $E' \in S(\beta')$ such that

$$|E - E'| < C|\beta - \beta'|^{1/2}. \quad (4.4)$$

In [5], the authors give a good bound on C requiring that $|\beta - \beta'|$ be sufficiently small. As the reader can verify, a trivial modification of the proof in [5] allows us to fix $C = 60$ for the almost Mathieu operator (worse than in [5]) but without any condition on β , $\beta' \in (0, 1)$. Thus we set $C = 60$.

This continuity property (4.4) for $\beta = p_n/q_n$, $\beta' = p_{n+1}/q_{n+1}$, together with the identity

$$\left| \frac{p_n}{q_n} - \frac{p_{n+1}}{q_{n+1}} \right| = \frac{1}{q_n q_{n+1}}$$

and the bound (4.1) implies that there exists $E' \in S(p_{n+1}/q_{n+1})$ such that

$$E' \in \left(\frac{E_2}{2} - \frac{C}{\sqrt{q_n q_{n+1}}}, \frac{E_2}{2} + \frac{C}{\sqrt{q_n q_{n+1}}} \right) \subset \left(\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^\kappa/2}}, \frac{E_2}{2} + \frac{C}{\sqrt{C_3 q_n^\kappa/2}} \right). \quad (4.5)$$

Using (4.2) and recalling that¹ $\kappa = 4C_1$, we see that

$$\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^\kappa/2}} > \frac{2}{C_2 q_n^{C_1}} - \frac{C}{\sqrt{C_3/2} q_n^{2C_1}} = \frac{2}{C_2 q_n^{C_1}} \left(1 - \frac{CC_2}{\sqrt{2C_3} q_n^{C_1}} \right),$$

¹ Note that here just $\kappa = 2C_1$ would do, cf a remark following Theorem 4.

and setting now

$$C_3 = 4^2 C^2 C_2^4 = 4^2 60^2 C_2^4, \quad (4.6)$$

we have

$$\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^\kappa/2}} > \frac{q_n^{-C_1}}{C_2}. \quad (4.7)$$

On the other hand, using (4.3), we obtain that

$$E_0 < \frac{8e}{C_3 q_n^{4C_1-1}} = \frac{e}{2C^2 C_2^4 q_n^{4C_1-1}} < \frac{q_n^{-C_1}}{C_2}.$$

Inequality (4.7) also shows that

$$\frac{E_2}{2} + \frac{C}{\sqrt{C_3 q_n^\kappa/2}} < E_2. \quad (4.8)$$

Thus,

$$E' \in (E_0, E_2),$$

which implies that

$$G_0(p_{n+1}/q_{n+1}) \subset (E_0, E_2 - \epsilon),$$

for some $\epsilon > 0$. The corresponding result for G_{-1} follows by the symmetry of the spectra. This proves the statement (a) of Theorem 4.

Now by the continuity (4.4) with $\beta = \alpha$, $\beta' = p_n/q_n$, and Theorem 3, we conclude that, for all $n = 1, 2, \dots$, there exists a gap $G_{n,2}(\alpha)$ of $S(\alpha)$ such that $G_{n,2}(\alpha) \cap G_0(p_n/q_n) \neq \emptyset$ and of length

$$\Delta_{n,2}(\alpha) > \Delta_0(p_n/q_n) - 2C|\alpha - p_n/q_n|^{1/2} > \frac{1}{C_2^2 q_n^{2C_1}} - \frac{2C}{C_3^{1/2} q_n^{\kappa/2}} = \frac{1}{2C_2^2 q_n^{\kappa/2}}. \quad (4.9)$$

We now verify that the gaps $G_{n,2}(\alpha)$, $G_{n+1,2}(\alpha)$ are distinct. Using the continuity once again, we obtain that there exists a point $E'' \in S(\alpha)$ such that

$$E'' \in \left(\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^\kappa}}, \frac{E_2}{2} + \frac{C}{\sqrt{C_3 q_n^\kappa}} \right). \quad (4.10)$$

Now it is easy to verify, similar to the calculations above, that

$$\frac{E_2}{2} - \frac{C}{\sqrt{C_3 q_n^\kappa}} > \frac{7}{4} \frac{1}{C_2 q_n^{C_1}}$$

and

$$E_0 + \frac{2C}{\sqrt{C_3 q_n^\kappa}} < \frac{1}{C_2 q_n^{C_1}},$$

and therefore (4.10) together with (4.8) yields

$$E'' \in \left(E_0 + \frac{2C}{\sqrt{C_3 q_{n+1}^\kappa}}, E_2 \right). \quad (4.11)$$

Thus $G_{n+1,2}(\alpha)$ lies to the left of E'' , and $G_{n,2}(\alpha)$ to the right of E'' , so that $G_{n,2}(\alpha)$ and $G_{n+1,2}(\alpha)$ are distinct gaps, $n = 1, 2, \dots$. Similar results for $G_{n,1}(\alpha)$ follow by the symmetry. This proves the statement (b) of Theorem 4.

The proof of the statement (c) is similar and based on (1.9). It is a simple exercise. \square

Acknowledgements. The work of the author was partially supported by the Leverhulme Trust research fellowship RF-2015-243. The author is grateful to Jean Downes and Ruedi Seiler for their hospitality at TU Berlin, where part of this work was written, and to the referees for very useful comments.

Open Access This article is distributed under the terms of the Creative Commons Attribution 4.0 International License (<http://creativecommons.org/licenses/by/4.0/>), which permits unrestricted use, distribution, and reproduction in any medium, provided you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made.

References

1. Aubry, S., André, G.: Analyticity breaking and Anderson localisation in incommensurate lattices. *Ann. Isr. Phys. Soc.* **3**, 133–164 (1980)
2. Avila, A., Krikorian, R.: Reducibility or nonuniform hyperbolicity for quasiperiodic Schrödinger cocycles. *Ann. Math.* **164**, 911–940 (2006)
3. Avila, A., Jitomirskaya, S.: The ten martini problem. *Ann. Math.* **170**, 303–342 (2009)
4. Azbel, M.Y.: Energy spectrum of a conduction electron in a magnetic field. *Sov. Phys. JETP* **19**, 634–645 (1964)
5. Avron, J., van Mouche, P., Simon, B.: On the measure of the spectrum for the almost Mathieu operator. *Commun. Math. Phys.* **132**, 103–118 (1990)
6. B ellissard, J., Simon, B.: Cantor spectrum for the almost Mathieu equation. *J. Func. Anal.* **48**, 408–419 (1982)
7. Choi, M.-D., Elliott, G.A., Yui, N.: Gauss polynomials and the rotation algebra. *Invent. Math.* **99**, 225–246 (1990)
8. Deift, P., Simon, B.: Almost periodic Schr odinger operators III. The absolutely continuous spectrum in one dimension. *Commun. Math. Phys.* **90**, 389–411 (1983)
9. Hardy, G.H., Littlewood, J.E.: Some problems of Diophantine approximation: a series of cosecants. *Bull. Calcutta Math. Soc.* **20**, 251–266 (1930). In: *Collected papers of G.H. Hardy* **1**, Oxford (1966)
10. Hardy, G.H., Littlewood, J.E.: Some problems of Diophantine approximation: the lattice points of a right-angled triangle (second memoir). *Abhandlungen Math. Sem. Hamburg. Uni.* **1**, 212–249 (1922). In: *Collected papers of G.H. Hardy* **1**, Oxford (1966).
11. Helffer, B., S j strand, J.: Semi-classical analysis for Harper’s equation. III. Cantor structure of the spectrum. *Mem. Soc. Math. France (N.S.)* **39**, 1–139 (1989)
12. Jitomirskaya, S., Ya., Krasovsky, I.V.: Continuity of the measure of the spectrum for discrete quasiperiodic operators. *Math. Res. Lett.* **9**, 413–422 (2002)
13. Kohmoto, M., Hatsugai, Y.: Peierls stabilization of magnetic-flux states of two-dimensional lattice electrons. *Phys. Rev. B* **41**, 9527–9529 (1990)
14. Korotyaev, E., Krasovsky, I.: Spectral estimates for periodic Jacobi matrices. *Commun. Math. Phys.* **234**, 512–532 (2003)
15. Krasovsky, I.V.: Bethe ansatz for the Harper equation: solution for a small commensurability parameter. *Phys. Rev. B* **59**, 322–328 (1999)
16. Last, Y.: Zero measure spectrum for the almost Mathieu operator. *Commun. Math. Phys.* **164**, 421–432 (1994)
17. Last, Y.: On the measure of gaps and spectra for discrete 1D Schr odinger operators. *Commun. Math. Phys.* **149**, 347–360 (1992)
18. Last, Y., Wilkinson, M.: A sum rule for the dispersion relations of the rational Harper’s equation. *J. Phys. A* **25**, 6123–6133 (1992)

19. Last, Y., Shamis, M.: Zero Hausdorff dimension spectrum for the almost Mathieu operator (**to appear in Commun. Math. Phys.**) [arXiv:1510.07651](https://arxiv.org/abs/1510.07651)
20. Liu, W., Yuan, X.: Spectral gaps of almost Mathieu operator in exponential regime. *J. Fractal Geom.* **2**(1), 1–5 (2015) [arXiv:1311.0658](https://arxiv.org/abs/1311.0658)
21. Mandelshtam, V.A., Zhitomirskaya, S.Y.: 1D-quasiperiodic operators. Latent symmetries. *Commun. Math. Phys.* **139**, 589–604 (1991)
22. van Mouche, P.: The coexistence problem for the discrete Mathieu operator. *Commun. Math. Phys.* **122**, 23–33 (1989)
23. Puig, J.: Cantor spectrum for the almost Mathieu operator. *Commun. Math. Phys.* **244**, 297–309 (2004)
24. Spencer, D.C.: On a Hardy–Littlewood problem of Diophantine approximation. *Proc. Camb. Phil. Soc.* **35**, 527–547 (1939)
25. Wiegmann, P.B., Zabrodin, A.V.: Quantum group and magnetic translations Bethe ansatz for the Azbel–Hofstadter problem. *Nucl. Phys. B* **422**, 495–514 (1994)

Communicated by J. Marklof